

# Counting Particles: a simple and fast surface reconstruction method for particle-based fluids

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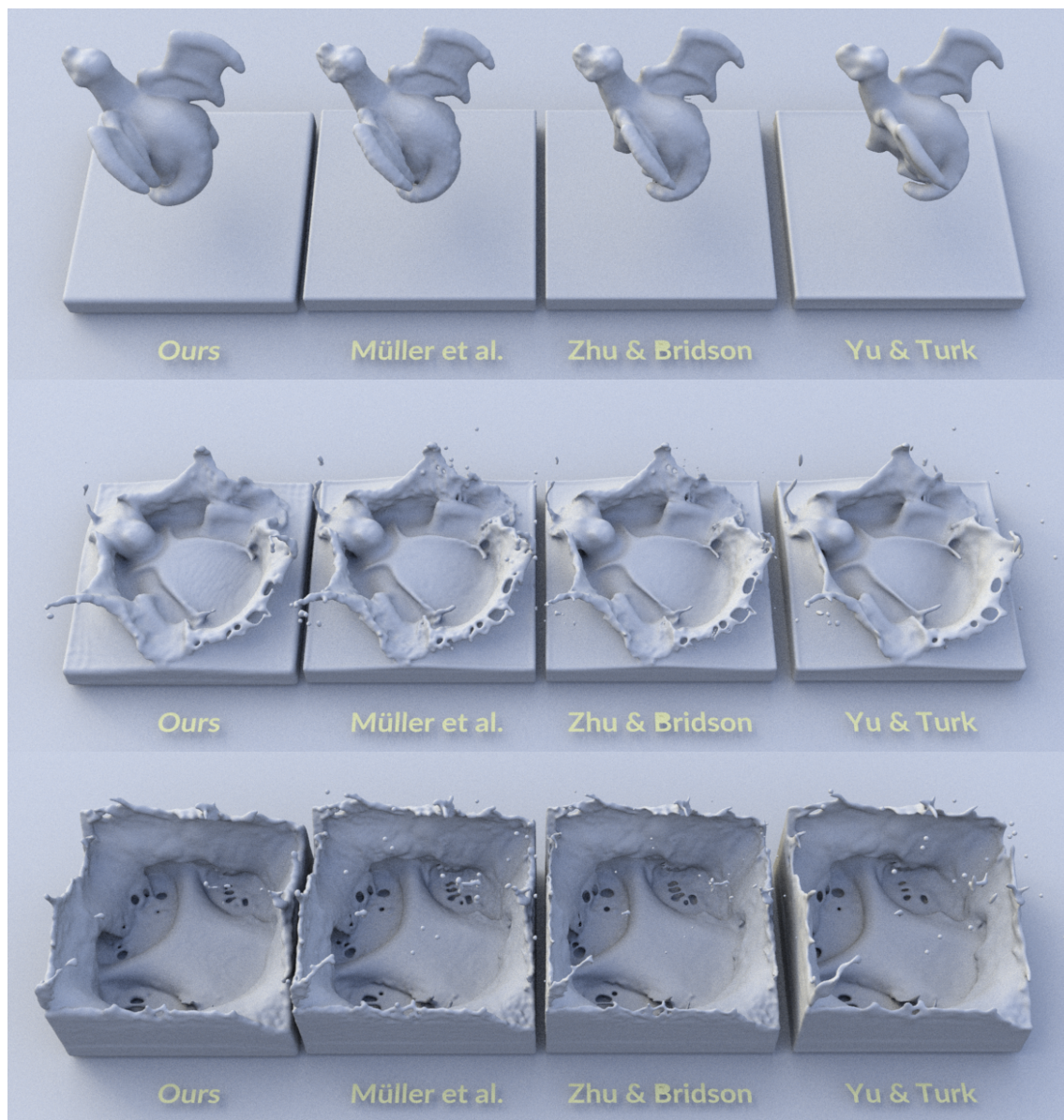


Fig. 1: Surface reconstruction of a liquid toy dragon with 1.2M particles in a grid with  $256^3$  cells: from left to right, our method, Müller et al. [1], Zhu and Bridson [2], Yu and Turk [3]. Our method is  $2.2\times$  faster than the best competitor method.

**Abstract**—We present a novel and efficient surface reconstruction for particle-based fluid methods. Although particle-based methods are practical for computing level-sets that represent liquid interfaces, these methods are computationally expensive when the number of particles increases considerably due to the intense usage of particle approximations. This paper introduces a

simple level-set approximation using a discrete indicator function (DIF) defined by counting particles inside grid cells. Our method is fast, easy to code, and can be adapted straightforwardly in particle-based solvers, even implemented in GPU. Moreover, we show the effectiveness of our approach through a set of experiments against prior surface reconstruction methods.

## I. INTRODUCTION

In particle-based fluid methods, such as Smoothed Particle Hydrodynamics (SPH) [4] and Position-Based Dynamics (PBD) [5], particles are usually employed to track the air-liquid interface in Lagrangian fluid flow simulations. Although these methods have been successfully used in interactive applications, rendering complex liquid animations with a high number of particles at reasonable computational times, even in GPU architectures, remains a subject of intense research in computer animation.

The rendering of air-liquid interfaces in particle-based fluids consists of two stages: first, splatting the level-set function defined by the particles to a regular grid. Second, the liquid surface reconstruction is given by the isosurface extracted from a discrete level-set using a polygonization algorithm, such as Marching Cubes (MC) [6], where a polygonal mesh represents the isosurface. However, this entire process is computationally expensive since the surface’s smoothness and topological correctness require a high-resolution grid.

On the other hand, in computational fluid dynamics, an indicator function  $\mathbb{1}_F$  (i.e., a function that is one on the interior of the fluid body  $F$  and zero on the exterior) is usually employed for tracking liquid interfaces [7]. Let  $\mathcal{G}$  be a regular grid of resolution  $n_x \times n_y \times n_z$  grid that encloses the domain. The key idea is to compute a local fraction of volume occupied by the fluid in each grid cell  $K \in \mathcal{G}$ , as follows:

$$\vartheta(K) = \frac{1}{V_K} \int_K \mathbb{1}_F(\mathbf{x}) d\mathbf{x}, \quad (1)$$

where  $V_K$  is the volume of  $K$ . The scalar field  $\vartheta : \mathcal{G} \rightarrow [0, 1]$  is known as *discrete indicator function* (DIF). Besides, there are robust surface reconstruction methods from DIFs in the literature [8], [9].

This paper presents a novel and practical surface reconstruction for particle-based fluid methods. We approximate efficiently a DIF by simply counting particles inside grid cells. Therefore, we show the effectiveness of our approach through a set of comparisons against prior surface reconstruction methods. Figure 1 shows our method in action.

In summary, the contributions of our method are:

- a novel DIF that uses only the number of particles inside the grid cells;
- our framework speeds up considerably the surface reconstruction compared with prior methods;
- our method is simple and easy to code, even in GPU.

## II. RELATED WORK

To better contextualize our method and highlight its properties we organize the existing methods for particle-based fluid rendering into two main groups: *mesh-based* and *screen-space* methods.

**Mesh-based methods.** The main goal of these methods is to extract a smooth triangle mesh from the particle positions using MC-based algorithms. Typically, the liquid surface is represented implicitly by the zero level-set of a signed distance

field computed from a weighted sum of kernel evaluations from the particles’ distances. These methods can use isotropic kernels [1], [2], adaptive size kernels [10] or anisotropic SPH kernels [3]. Despite the existence of parallel implementations of these methods [11]–[13], if the liquid spreads more over the computational domain, the underlying MC grid and its resulting surface mesh become very large, causing excessive memory consumption. Bhattacharya et al. [14] improved the undesired blobby appearance of the level-sets using a smoothing process by solving a constrained optimization problem. Sandim et al. [15] proposed an alternative framework for surface reconstruction. Their framework relies on a level-set definition using the Hermite data (particle positions and normals) from the boundary particles. The liquid surface is obtained fitting the boundary particles using Screened Poisson surface reconstruction [16]. However, this method also suffers the same issues of the kernel-based methods.

**Screen-space methods.** This class of methods performs in 2D image space using a smoothed depth buffer from the visible liquid surface defined by spherical particles, where the resulting surface is represented without mesh generation by using rasterization techniques [17]–[21]. The liquid surface’s visual quality relies on the depth buffer’s image-based filtering process, which may demand large convolution kernels and perform multiple filter iterations. However, beyond the screen-space size and the number of filter passes, these methods’ efficiency also depends on the number of particles. Recently, Oliveira and Paiva [22] improved the prior screen-space methods computing volumetric rendering effects in a small subset of particles located at a narrow-band of the air-liquid interface.

Despite the proposed approach belonging to the class of mesh-based methods, our method was strongly influenced by screen-space methods, extending the filtering process to a 3D image (a DIF in our case).

## III. THE METHOD

In this section, we explain the pipeline of the proposed surface reconstruction method. Given an input particle system  $\mathcal{P}^t$  at time-step  $t$ , our method performs three main steps, as illustrated by Figure 2.

**DIF evaluation.** In this step, we approximate the local volume of fluid simply by counting the particles inside each cell  $K \in \mathcal{G}$ . Let  $N_K^t$  be the number of particles of  $\mathcal{P}^t$  inside  $K$ . Firstly, we compute the initial particle average  $\mu_0$  (at time-step  $t = 0$ ) given by:

$$\mu_0 = \frac{1}{|\mathcal{F}|} \sum_{K \in \mathcal{F}} N_K^0,$$

where the operator  $|\cdot|$  denotes the set’s cardinality and  $\mathcal{F} \subseteq \mathcal{G}$  is the subset of *full cells* of  $\mathcal{G}$ , i.e., formed by cells that contain particles in their interior. Assuming that the cell volume  $V_K$  is entirely occupied by the volume of  $\mu_0$  particles, i.e.,  $V_K = \mu_0 V_p$ , where  $V_p$  is the particle volume. Thus, we approximate the Equation (1) as follows:

$$\tilde{\vartheta}(K) = \frac{N_K^t V_p}{V_K} = \frac{N_K^t V_p}{\mu_0 V_p} = \frac{N_K^t}{\mu_0}.$$

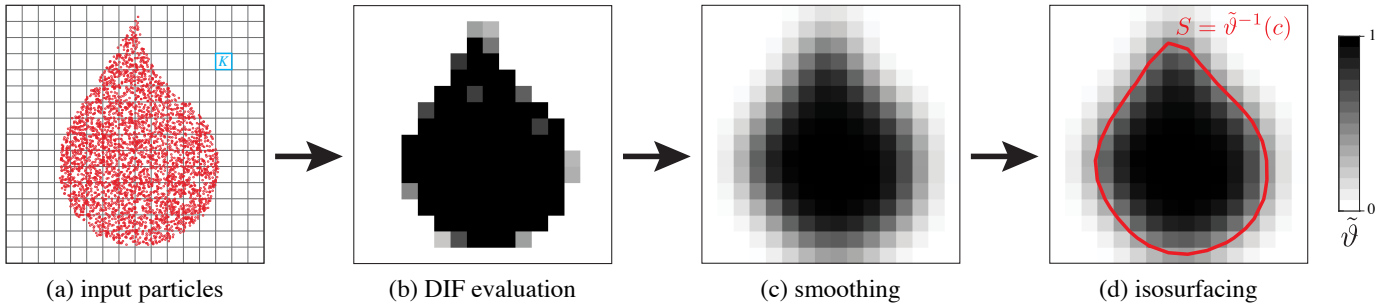


Fig. 2: Surface reconstruction pipeline.

In order to ensure  $\tilde{\vartheta}(K) \in [0, 1]$ , the fluid occupancy in a cell  $K$  in our approach is given by:

$$\tilde{\vartheta}(K) = \frac{\min(N_K^t, \mu_0)}{\mu_0}. \quad (2)$$

**Smoothing.** To reduce the DIF discontinuities around the liquid interface, we smooth the field (2) by applying 3D blur filters from image processing [23], e.g., box and Gaussian filters. Firstly, we apply a box filter of size  $3 \times 3 \times 3$  to eliminate small holes (i.e., empty cells surrounded by full cells). Then, we apply Gaussian filter of size  $5 \times 5 \times 5$  with standard deviation  $\sigma = 1.2$  in a single pass to enhance the DIF.

**Isosurfacing.** In our method, given a cell  $K_i \in \mathcal{G}$ , we assume that the smoothed DIF  $\tilde{\vartheta}$  is sampled at cell centers  $\mathbf{p}_i$  of  $K_i$ . The liquid surface is represented by the level-set  $S = \tilde{\vartheta}^{-1}(c)$  with isovalue  $c \in (0, 1)$ . Thus, once computed the field  $\tilde{\vartheta}$ , we have fractional volumes of fluid  $\tilde{\vartheta}_i$  located at the centers  $\mathbf{p}_i$ . Then, we extract the polygonal mesh of the level-set  $S$  executing the MC algorithm in the *dual grid*. The cells of the dual grid are obtained by connecting the centers of the adjacent cells of  $\mathcal{G}$  (see Figure 3). The MC’s lookup table determines the local topology of  $S$  inside each dual cell by indexing the configurations of  $\text{sign}(\tilde{\vartheta}_i - c)$  at the eight corners of the cell. Considering a linear approximation of  $\tilde{\vartheta}$ , given a dual edge  $e_{ij} = (\mathbf{p}_i, \mathbf{p}_j)$  where  $(\tilde{\vartheta}_i - c) \cdot (\tilde{\vartheta}_j - c) < 0$ , the intersection point (vertex)  $\mathbf{p}$  of  $S$  with  $e_{ij}$  is computed as follows:

$$\mathbf{p} = (1 - \alpha)\mathbf{p}_i + \alpha\mathbf{p}_j \quad \text{with} \quad \alpha = \frac{c - \tilde{\vartheta}_i}{\tilde{\vartheta}_j - \tilde{\vartheta}_i}.$$

After processing each dual cell, the entire surface  $S$  is extracted. In our experiments, we use the isovalue  $c = 0.25$ .

#### IV. RESULTS

We implemented our approach in C++ and a parallel version of our code on GPU using CUDA. The particle-based fluid simulations were produced with SPH using the computational framework provided by DualSPHysics [24]. All results have been achieved using a computer with AMD Ryzen 9 3950X and 32GB RAM and NVIDIA GeForce RTX 2070 with 8GB VRAM.

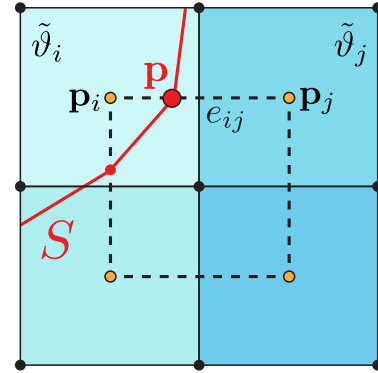


Fig. 3: Isosurfacing of a DIF in 2D: a dual grid is obtained from the grid  $\mathcal{G}$  (solid black lines). For each dual cell (dashed black line), MC examines the configuration of  $\text{sign}(\tilde{\vartheta}_i - c)$  at the corners  $\mathbf{p}_i$  (orange dots) to define the local topology of the surface  $S$  (solid red line) and determines the intersection points  $\mathbf{p}$  (red dots) between  $S$  and the dual edges  $e_{ij}$ .

Figures 1, 4, and 5 show comparisons of our surface reconstruction method applied in different simulations in comparison with previous methods proposed by Müller et al. [1], Zhu and Bridson [2], and Yu and Turk [3]. Furthermore, the implementations in C++ of these methods can be found in the Github<sup>1</sup> repository from Kim’s book [25].

Table I shows the computational times and some statistics for a set of experiments presented in this section. The column  $|\mathcal{P}|$  is the number of particles and the column **res** is the resolution of  $\mathcal{G}$ . The average **time** of each method across all animation frames was measured using a single-core CPU. The *speedup* (in parenthesis) shows how fast our approach is compared to prior methods. Note that our approach is  $2.1\times$  faster in the worst case and  $13.9\times$  faster in the best case, demonstrating our method’s efficiency.

#### V. DISCUSSION

**Scalability and profiling.** Figure 6 shows the computational timing of our approach implemented on GPU and the performance profiling of each stage of the reconstruction pipeline (as shown in Figure 2) with different resolutions of  $\mathcal{G}$ . As can

<sup>1</sup><https://github.com/doyubkim/fluid-engine-dev>

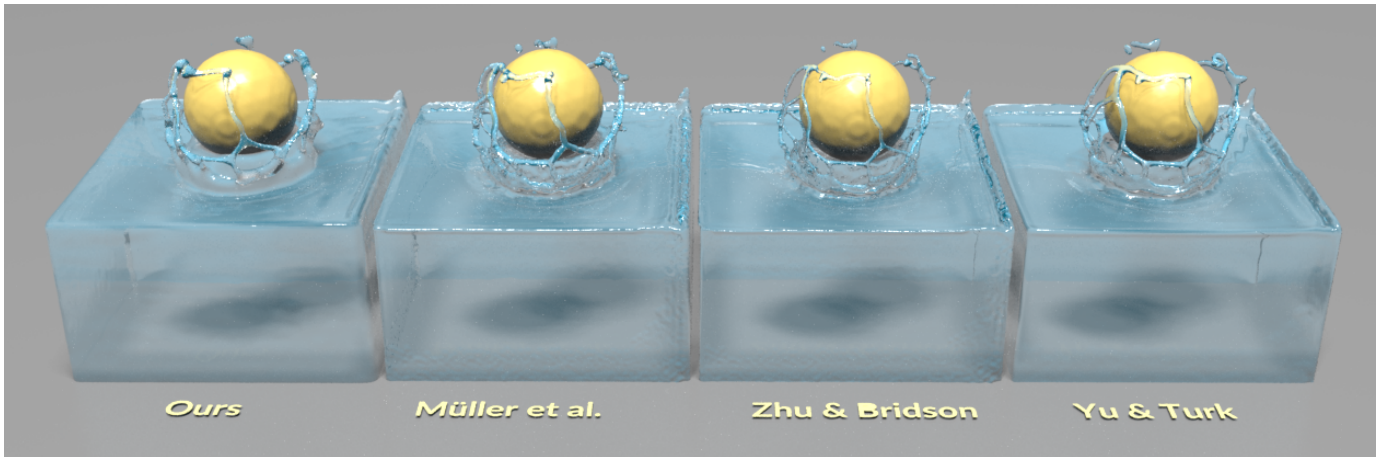


Fig. 4: Surface reconstruction of a floating ball with 0.93M SPH particles in a grid with  $256^3$  cells: from left to right, our method, Müller et al. [1], Zhu and Bridson [2], Yu and Turk [3].

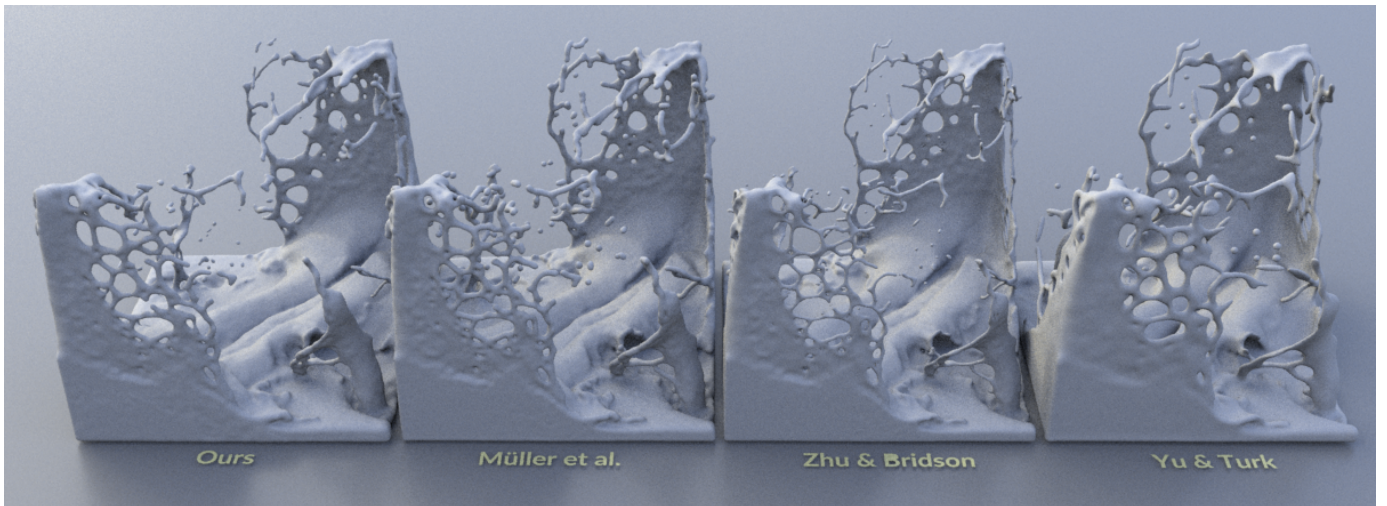


Fig. 5: Surface reconstruction of a double dam-break simulation with 1M SPH particles in a grid with  $256^3$  cells: from left to right, our method, Müller et al. [1], Zhu and Bridson [2], Yu and Turk [3].

TABLE I: Average statistics and computational times (in seconds) per frame.

Experiment	$\mathcal{P}$	res	time (speedup)			
			Müller et al. [1]	Zhu and Bridson [2]	Yu and Turk [3]	Ours
Toy dragon (Fig. 1)	1.20M	$256^3$	22.42 (3.0)	15.94 (2.2)	85.03 (11.6)	7.30 (-)
Floating ball (Fig. 4)	0.93M	$256^3$	18.85 (2.6)	16.23 (2.3)	100.05 (13.9)	7.21 (-)
Double dam-break (Fig. 5)	1.00M	$256^3$	19.48 (2.9)	14.05 (2.1)	76.27 (11.3)	6.74 (-)

be seen, the computational time related to grid operations to produce the smoothed DIF increases when we refine the grid, becoming a potential bottleneck in high-resolutions. Regarding the rendering, our method preserves nicely small-scale liquid details even in a grid with a resolution of  $512^3$ . Important to note that the GPU version is almost  $100\times$  faster than the single-core version for a grid resolution of  $256^3$  (see Table I).

**Limitations and future work.** The grid representation restricts our GPU implementation to bounded domains. It opens possibilities for replacing our current data structure

with sparse grid representations using GVDB Voxels [26]. Aggressive smoothing can remove small surface details like liquid droplets. Thus, another direction of future research is constructing a “detail-aware” blur using adaptive filters [27] for DIF smoothing and more sophisticated polygonization algorithms suited for DIFs [8], [9] as well.

## VI. CONCLUSION

We introduced a simple and fast surface reconstruction method for liquid interfaces suited for particle-based fluid solvers on both CPU and GPU architectures. The proposed

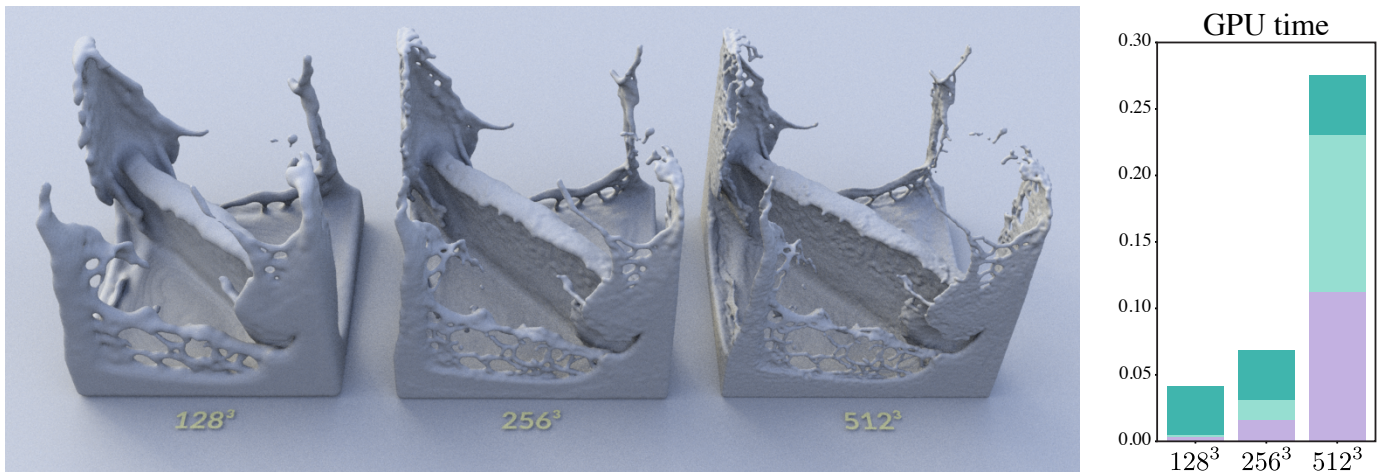


Fig. 6: Analysis of the grid resolution on GPU: surface reconstruction of the liquid splashing in the double dam-break (Figure 5) using our method with different grid resolutions (left) and the average computational timing (in seconds) of each pipeline stage (right): DIF evaluation (■), smoothing (■), and isosurfacing (■).

method relies on a novel smoothed DIF defined by counting particles inside grid cells, providing a high-quality surface. Our method provides a significant speed-up for surface reconstruction compared to the prior methods, as attested by the set of experiments and comparisons carried out in the paper.

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